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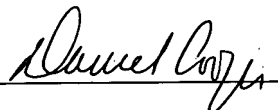
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METHOD FOR IMPROVING DEPTH DISCRIMINATION IN OPTICAL REPRODUCTION SYSTEMS

[0001] The invention is directed to a method for depth discrimination in optically imaging systems. In particular, it is applicable in light microscopy for achieving an improvement in image quality in the examination of three-dimensionally extending objects.

[0002] A problem arising in conventional light microscopy when examining three-dimensionally extending objects, i.e., objects whose extension along the optical axis is greater than the depth of focus of the objectives that are used, consists in that the in-focus image from the region of focus is superposed with extra-focal image components which are accordingly imaged out of focus.

[0003] In order to overcome this problem, it is known to use confocal imaging in which the light originating from the area outside the region of focus is excluded by means of a pinhole and therefore does not contribute to the image. In this way, an optical section is formed. This confocal point-type imaging requires that the object is scanned in the image plane in order to obtain an image. This scanning can be carried out either by means of scanners in laser scanning microscopes or by means of Nipkow disks.

[0004] By recording a plurality of optical section images in different focus positions, a z-stack can be obtained so that the object may be displayed three-dimensionally.

[0005] Another method for generating optical sections is to use structured illumination. This was first stated by Meir Ben-Levy and Eyal Pelac in WO 97/6509. This method was improved and expanded as described by Tony Wilson et al. in WO 98/45745 and by Volker Gerstner et al. in WO 02/12945. The disclosure of these three publications is expressly referenced herein.

[0006] In WO 97/6509, the object is illuminated by a periodic structure (sine grating or rectangular grating) and an image of the object is recorded by a camera, digitized and stored in a storage. The periodic structure is subsequently displaced within the image plane in such a way that the phase position of the structure is changed and an image is again recorded and stored. This process (displacement, image recording, storage) can be repeated many times.

A section image is then generated by calculating from the existing images. The indicated mathematical formulation is a Fourier expansion, which leads to a complicated formula apparatus.

[0007] WO 98/45745 shows a simpler formula for the section images which can be derived by simplifying the formula from WO 97/6509 in case of equal phase displacements between the individual recordings.

[0008] Realizing the physical boundary conditions required for the application of the indicated method proves very difficult in practice. For example, variation in lamp brightness between the different recordings leads to stripe artifacts in the generated section images. With fluorescing objects, additional problems occur due to the time-dependent fading of the fluorescent dyes, which likewise results in errors. The necessary constancy of the individual phase displacement steps cannot be maintained in practice.

[0009] Therefore, it was suggested in WO 02/12945 to compensate for the influence of lamp brightness that varies over time by coupling out part of the light serving to illuminate the object, registering the intensity and subsequently scaling the individual recordings. In order to take into account unequal phase displacement steps, a system of equations (Equation 22, op. cit.) is indicated. To compensate for fading or bleaching, instead of the necessary minimum of three recordings per section image, it is suggested that six recordings are registered in the sequence 1 - 2 - 3 - 3 - 2 - 1, two section images are calculated (from 1 - 2 - 3 and 3 - 2 - 1) and the average is determined therefrom.

[0010] A considerable expenditure on instrumentation is required in order to realize these suggestions. Further, the recording of additional images prolongs the required recording time and therefore also increases loading of the sample by illumination with the fluorescence excitation light.

[0011] It is the object of the invention to overcome the disadvantages of the prior art and to provide an improved method for determining optical section images. This object is met by a method for improving the depth discrimination of optically imaging systems according to the main claim. Advantageous developments are indicated in the subclaims.

[0012] The invention will be described more fully in the following with reference to the drawings.

[0013] Fig. 1 shows the basic construction of a microscope with structured illumination;

[0014] Fig. 2 is a schematic view of the bandpass filter, according to the invention, in the Fourier plane;

[0015] Fig. 3 is a schematic view of the position of the images of the periodic structure;

[0016] Fig. 4 shows a flowchart for determining correction values through linear optimization.

[0017] Fig. 1 is a simplified optical schematic illustrating the structured illumination with transmitted illumination by way of example. The imaging beam path (image-forming beam path) is shown. A one-dimensional periodic structure (transmission grating) (3) which is located in a focus plane of the optical arrangement shown in the drawing is illuminated by a light source (1) through collector optics (2). The grating is followed in light direction by a plane-parallel glass plate (4). The angle of the plane-parallel plate relative to the optical axis can be adjusted in a defined manner. The structure is imaged in the specimen plane (7) by the illumination-side optics (5, 6) (condenser). Imaging is effected by the light coming from the specimen through a pair of lenses (8, 9) (objective and tube lens) in the focus plane (10) following the latter in which, e.g., a CCD matrix of a digital camera can be arranged. By tilting it in a defined manner, the plane-parallel glass plate (4) serves to displace the image of the grating structure (3) on the object in the specimen plane (7). For incident fluorescence observation, the objective (8) preferably serves at the same time as a condenser. As was described in WO 02/12945, the resulting brightness distribution in the object is registered in at least 3 positions of the glass plate (4) by means of the digital camera. In this connection, the brightness distribution $I_i(x, y)$ for a sine/cosine grating can be stated in a simplified manner by the following formula:

$$I_i(x, y) = I_0(x, y) \cdot (1 + m(x, y) \cdot \cos(\phi_0(x, y) + \phi_i(x, y))), \quad (1)$$

where $i=0 \dots N-1$ is the i th phase position of the projected grating, and N is the quantity of recordings, $m(x, y)$ is the modulation depth of the object (and therefore the image information sought at point x, y), and ϕ_i are the phase values. This equation contains the

three unknown variables I_0 , m and ϕ_0 . Accordingly, these variables can be determined by means of at least three measurements with specifically varied ϕ_i ($i = 1, 2, 3$). The solution can be found from the measurements through a least square formulation. For this purpose, it is possible to represent Equation (1) in a more compact form and to rewrite the cosine function using the addition theorem:

$$I(x, y, \phi_i) = a_0(x, y) + a_1(x, y) \cdot f_1(\phi_i) + a_2(x, y) \cdot f_2(\phi_i), \quad (2)$$

where

$$f_1(\phi_i) = \cos \phi_i$$

$$f_2(\phi_i) = \sin \phi_i$$

$$a_0(x, y) = I_0(x, y)$$

$$a_1(x, y) = I_0(x, y) \cdot m(x, y) \cdot \cos \phi_0(x, y)$$

$$a_2(x, y) = -I_0(x, y) \cdot m(x, y) \cdot \sin \phi_0(x, y)$$

[0018] The functions f_1 and f_2 therefore depend only on the phase displacements ϕ_i , which can be freely selected in principle. For a structure that is not sinusoidal but is periodic, the brightness distribution $I_i(x, y)$ can also be approximated by series expansion. The principle of calculation remains basically the same. Expressed in matrix form, the least square solution is:

$$\hat{M} \cdot \vec{a} = b, \quad (3)$$

where

$$\hat{M} = \begin{pmatrix} N & \sum_N f_1(\phi_i) & \sum_N f_2(\phi_i) \\ \sum_N f_1(\phi_i) & \sum_N f_1^2(\phi_i) & \sum_N f_1(\phi_i) f_2(\phi_i) \\ \sum_N f_2(\phi_i) & \sum_N f_1(\phi_i) f_2(\phi_i) & \sum_N f_2^2(\phi_i) \end{pmatrix}, \quad (4)$$

where N is the quantity of measurements (in this case, phase steps) and

$$\vec{a} = \begin{pmatrix} a_0(x, y) \\ a_1(x, y) \\ a_2(x, y) \end{pmatrix} \quad (5)$$

and

$$\vec{b} = \begin{pmatrix} \sum_N I(x, y, \phi_i) \\ \sum_N [I(x, y, \phi_i) \cdot f_1(\phi_i)] \\ \sum_N [I(x, y, \phi_i) \cdot f_2(\phi_i)] \end{pmatrix} \quad (6)$$

[0019] The solution to this matrix equation is obtained mathematically by inverting the matrix M.

$$\begin{pmatrix} a_0(x, y) \\ a_1(x, y) \\ a_2(x, y) \end{pmatrix} = \hat{M}^{-1} \cdot \vec{b} \quad (7)$$

[0020] Accordingly, the unknown variables being sought can be determined from the measurements, represented by vector b, and the phase displacements of the grating. In particular, the modulation depth at point (x, y) is given by:

$$m(x, y) = \frac{\sqrt{a_1^2(x, y) + a_2^2(x, y)}}{a_0(x, y)} \quad (8)$$

[0021] Since the grating contrast and, therefore, by necessity, the modulation depth in the object decreases with defocusing of the object, Equation (8) is a conditional equation for obtaining a depth-discriminated image (optical section). Equations (2) and (3) provide a universal approach that allows the phase positions ϕ_i to be adapted to the actual conditions. Accordingly, the number of steps and the positions in the object space (location of the imaged grating lines) can be freely selected.

[0022] Each optical system, according to its optical transfer function, transfers spatial frequencies up to a threshold value that is determined by the numerical aperture NA of the optics and the light wavelength λ :

$$k_{\max} = \frac{4\pi NA}{\lambda}$$

[0023] As a result, the higher harmonics of the grating period of a grating which is imaged in the object space and which is not a pure sine grating are transmitted up to this threshold frequency k_{\max} and are found in the Fourier-transformed image of the object as local maxima. These components of the grating are also contained in a section image determined according to the plan detailed above and result in stripe artifacts. These artifacts can be eliminated in that the harmonics of the projected grating are localized in the spatial frequency spectrum by Fourier transformation and are removed by a specific setting of filters (bandpass filters, bandstops). A subsequent inverse transformation then results in images that are free of artifacts. This method is shown schematically in Fig. 2. In the Fourier plane (11), the higher harmonics of the grating structure (3) are designated by (12), the zero point is designated by (13) and corresponds to the DC component of the illumination, i.e., a uniform, unstructured illumination. By introducing bandpass filters corresponding to the harmonics (12) into the beam path, these areas do not contribute to the imaging and are therefore not registered by the digital camera.

[0024] With respect to fluorescing objects, there is another reason for the occurrence of stripe artifacts, namely, fluorescence bleaching (photobleaching, fluorescence fading). In this case, the imaging of the structure (e.g., of the grating) in the object results in regions of differing degrees of bleaching of the fluorescence and therefore, finally, in stripe artifacts in every image recorded subsequently. The frequency components of the artifacts in the spectrum of an image of this type are particularly strong in the fundamental frequency of the projected structure. By removing the spatial frequency components of the projected structure in the Fourier-transformed image by means of corresponding bandpass filters and subsequent inverse transformation, an artifact-free image can also be generated in this case in principle.

[0025] Prevention of these artifacts occurring in fluorescence illumination can also be achieved by determining and making use of the bleaching characteristics of fluorescing objects in the sequential recording of images of the same object. In this case, the illumination results in bleaching of the object at the illuminated points for every image recording. Apart from the illumination intensity, the degree of bleaching also depends on the photochemical and photophysical characteristics of the fluorescing object. By determining the fluorescence intensity of the object through integration of the intensities over all object points or selected object points (x, y), it is possible to rescale the image sequence after the images have been

recorded. For this purpose, the intensities at the image points (x, y) are scaled to a maximum value either by means of an analog camera and analog-to-digital converter or a CCD camera using image information stored in a computer by determining the quotients of the fluorescence intensities of the images by pairs. In this way, the fluorescence loss due to fluorescence bleaching is taken into account. In addition, according to Fig. 3, regions of interest (ROI) can be defined which are equally illuminated in both image recordings and which subsequently fluoresce; the degree of the reduction in fluorescence in these regions is a measurement for the bleaching rate and can therefore be used for scaling the different recordings of the brightness distribution. These ROIs can be defined by the user by means of corresponding input devices or can be defined automatically (from the respective position of the imaged structure). In Fig. 3, the ratios in a first recording of the brightness distribution on the object is shown at (14) and the ratios in a second recording with changed phase position of the imaged structure (16) are shown at (15). The regions of interest (17) are advantageously defined in such a way that uniform illumination is carried out in both recordings.

[0026] The method according to the invention allows optical section images to be determined practically for any selected phase displacements ϕ_i . In order to prevent errors, the phase displacements ϕ_i must be determined with high accuracy. This can be ensured in a simple manner by the calibration method described in the following, in which the influences of the entire system are taken into account.

[0027] For this purpose, a reference image (object with superimposed grating) is initially recorded and registered. A slight phase displacement of the grating to an initially unknown phase position of the grating is then carried out (this "displacement per control signal" is to be calibrated first) and a new image recording is made. The two images that are obtained are then compared by computation. Without limiting generality, this can be carried out by differentiating, summing, or in any other way that results in a merit function. The initial result is generally a striped image. The steps comprising displacement of the grating by a small amount, image recording, and comparison to the reference image are repeated until the merit function reaches the extreme value, that is, e.g., a differential image comprises only background noise or a sum image reaches maximum values. When the latter is accomplished, the calibration is terminated and a control value obtained in this way can be

kept in a storage medium so that it can be retrieved at a later time. Therefore, this value corresponds exactly to the displacement of the grating by a full period. Alternatively, it is also possible to assess a sum image for uniform brightness distribution (i.e., disappearance of stripe structure). In this case, this value corresponds to the displacement of the grating by half a period. This procedure is preferably carried out with a mirror as object. The described process can be carried out manually or automatically.

[0028] In the following, an alternative method for determining the influences of variations in lamp brightness, bleaching of the object in fluorescence illumination, and non-sinusoidal gratings is described with reference to Fig. 4.

[0029] For this purpose, the change in the illumination over time will be described by intensity-proportional scalar factors $\theta_i > 0$ which are taken into account in every successive recording (the dependency of the variables \mathbf{g}_i , \mathbf{o}_i , etc. on location (x, y) is no longer explicitly indicated in the following for the sake of simplicity):

$$\mathbf{g}_i = \theta_i \mathbf{o}_i \quad (9)$$

[0030] The modeled observation \mathbf{g}_i is the product of the ideal, consistent phase image \mathbf{o}_i . In order to describe a faulty system completely, only N-1 factors of θ_i need to be determined.

[0031] The bleaching of the object in fluorescence illumination depends on the effective radiation intensity over time. This varies spatially because of the projected grating and therefore must be handled in such a way that the grating function is taken into account. A simple description is possible when the bleaching is assumed to be a linear function of the illumination intensity:

$$\mathbf{k}_i = (1 - d \mathbf{o}_{i-1} \mathbf{k}_{i-1}) \prod_{j=0}^{i-1} \mathbf{k}_j, i \geq 1, \text{ where } \mathbf{k}_0 = \mathbf{1} \quad (10)$$

[0032] The vector $0 < \mathbf{k}_i \leq \mathbf{1}$ comprising spatially varying factors describes the attenuation as a function of the shape of the irradiation intensity and a selectable measurement $1 > d \geq 0$ for the bleaching when multiplied component-wise by the observations \mathbf{o}_i that have not yet undergone bleaching:

$$\mathbf{g}_i = \mathbf{o}_i \mathbf{k}_i \quad (11)$$

[0033] It follows from (10) that every observation is always multiplied by the factors of all preceding phase recordings; the latter are already bleached in part and overlap with the instantaneous phase. Alternatively, \mathbf{k}_i can also be described exponentially, i.e., corresponding to the decay curve of a determined fluorochrome.

[0034] Non-sinusoidal grating shapes cause stripe artifacts particularly for integral multiples of the grating frequency. With exact knowledge of frequency, phase position and curve shape of the grating in the individual recordings, it is possible to generate a sinusoidal shape. A correction vector \mathbf{I}_{Corr} similar to (10), but in this case along the grating periodicity, is considered in (12).

$$\mathbf{I}_{Corr} = \frac{f_{sin}}{f_{grating} + s} \quad (12)$$

[0035] The simulation of the observed grating function is characterized by $f_{grating}$. An ideally sinusoidal intensity curve is expressed by f_{sin} . The constant s in the denominator should take over the division by very small values of $f_{grating}$. The simulation of the observed grating function can be carried out in different ways. A simple possibility for approximation is provided, e.g., through the use of a Fourier series for the trapezoid function.

$$f_{grating}(x; b) = \frac{1}{2} + \frac{2}{b\pi} \sum_{i=1}^M \frac{\sin(2bi - b)}{(2i - 1)^2} \sin(2xi - x) \quad (13)$$

[0036] The edge steepness of the trapezoid transitions can be varied by b , x is the spatial coordinate, and M is the quantity of Fourier coefficients used. Alternatively, more exact models can also be used. They require convolution of the true grating function with a lateral component of the PSF of the microscope. A phase-synchronous synthesizing of the correction vector \mathbf{I}_{Corr} is crucial. When this is ensured, a corrected result for the observation can be achieved similar to (11) by a simple component-wise multiplication.

$$\mathbf{g}_{iCorr} = \mathbf{g}_i \mathbf{I}_{iCorr} \quad (14)$$

[0037] The corrected phase images are now obtained from:

$$I_i = \mathbf{o}_i = \frac{\mathbf{g}_i \mathbf{I}_{iCorr}}{\mathbf{k}_i \theta_i}, \text{ where } \theta_i > 0 \cap \mathbf{k}_i > 0 \quad (15)$$

By inserting (15) in Equations (3) or (8), a general formula is obtained for reconstruction which, moreover, contains the parameters θ_i , d and b .

[0038] The stated problem can be solved, for example, by converting to an extremal problem. The parameters θ_i , d and b can be determined with the aid of numeric optimization using a merit function.

$$M(\theta_i; d; b) = \alpha_0 |F\{\bar{a}\}_0|^2 + \alpha_i |F\{\bar{a}\}_\omega|^2 + \alpha_2 |F\{\bar{a}\}_{2\omega}|^2 + \dots + \alpha_n |F\{\bar{a}\}_{n\omega}|^2 \rightarrow \min \quad (16)$$

In this case, $F\{\bar{a}\}_\omega$ is a component of the Fourier transformation of the resultant vector \bar{a} from Equation (5). Functional transformations other than the Fourier transformation are also suitable in this connection.

[0039] The merit function (Equation (16)) is mentioned here only by way of example. Other merit functions are also applicable. The aim is to vary θ_i , d and b in such a way that (16) is minimized. Many methods are available for this purpose which are by no means limited to gradient methods or line-search methods. Further, every transformation coefficient is weighted by a coefficient α_i . Accordingly, the algorithm can be adapted to different signal-to-noise distances or preferred frequencies. An advantageous value in this respect for weighting the DC component is given by the following equation:

$$\alpha_0 \propto \frac{|F\{\bar{a}\}_1|^2 + |F\{\bar{a}\}_2|^2 + \dots + |F\{\bar{a}\}_n|^2}{|F\{\bar{a}\}_0|^2} \quad (17)$$

[0040] The general operation of the correction apparatus described herein is illustrated by the schematic flowchart in Fig. 4. The use of projections is advantageous for efficient determination of parameters θ_i , d and b . However, the determination of parameters is not limited to this one-dimensional case, but functions in the plane in an analogous manner.

[0041] With a fixed focus position of the camera, at least three brightness distributions (18), (19), (20), and possibly more (21), are registered in different phase positions $i=1, 2, 3, \dots, N$ of the grating structure imaged on the object. These two-dimensional brightness distributions are converted, respectively, to a one-dimensional distribution (22), (23), (24), and possibly (25), by means of projection. With initial values for the parameters θ_i , d and b ,

corrected phase images are calculated according to formula (15) and a first image of the optical section is calculated (26) therefrom by means of the formulation from formula (3). The initial values of parameters θ_i , d and b can be entered by the user or also estimated automatically. Known minimizing or maximizing tools of a merit function can be used for this estimation (e.g., Simplex, Marquardt, or the like). The merit function according to formula (16) is determined from this result in (27) and a test is made in (28) to determine whether or not a minimum has been reached. If this minimum has not yet been reached, a new estimation of the parameters θ_i , d and b is carried out (29) and continued with (26). When the minimum of the merit function has been reached, i.e., the value of parameters θ_i , d and b can be increased again by further variation of parameters θ_i , d and b , the optimal parameters are determined and the optical section image is determined in (30) by means of formulas (15), (3) and (8) and outputted in (31). Subsequently, the focus position can be changed and another optical section image can be determined, and so on, by means of the described method. By means of this scanning of the object in z-direction, a z-stack of section images and, accordingly, a three-dimensional representation of the object can be generated.

[0042] Another approach for taking into account bleaching effects in the object in fluorescence illumination is based on determination of a local correction function. For this purpose, a correction function $\kappa_i(x)$ is assumed:

$$\kappa_i(x) = \frac{\int_{x-\tau/2}^{x+\tau/2} g_i(\xi) d\xi}{\int_{x-\tau/2}^{x+\tau/2} g_i(\xi) d\xi}, \quad (18)$$

whose local value is formed by a grating period $\tau = 2\pi/\omega$. The formulation of this function can also take into account the dependency on y, although this is omitted in the present case for the sake of simplicity. Using a spatially variable bleaching function $\theta_i(x)$ and by suitable conversions with the approximation (see also formulas (1) and (9)

$g_i(x) = \theta_i(x)I_0(x)[1 + m \cos(\omega x + \varphi_i)]$, it can be shown by the following equation

$$\kappa_i(x) = \frac{\int_{x-\tau/2}^{x+\tau/2} I_0(\xi) d\xi + m \int_{x-\tau/2}^{x+\tau/2} I_0(\xi) \cos(\omega \xi + \varphi_1) d\xi}{\int_{x-\tau/2}^{x+\tau/2} \theta_i(\xi) I_0(\xi) d\xi + m \int_{x-\tau/2}^{x+\tau/2} \theta_i(\xi) I_0(\xi) \cos(\omega \xi + \varphi_1) d\xi}$$

that this correction function is proportional to $1/\theta_i(x)$ with uniform bleaching function.

[0043] It has proven advantageous to average over more than one grating period. Further, it has been shown that this correction function tends to spike in the neighborhood of edges in the bleaching function and can therefore generate additional errors in the results. In order to remedy this, it has proven successful to detect the occurrence of spiking by comparing to a mean value and to use the global value $1/\theta_i$ previously determined from formulas (9) and (16) at these points instead of the falsified value. The corresponding formula is then given by

$$\kappa_i(x) = \begin{cases} \kappa_i(x) & \text{for } \left| \kappa_i(x) - \int_{x-\tau}^{x+\tau} \kappa_i(\xi) d\xi \right| < \varepsilon \\ 1/\theta_i & \text{otherwise} \end{cases} \quad (19)$$

[0044] For this purpose, the threshold ε is advantageously set as a percentage of the variation of the local variability of the correction function within the image, e.g., 5%. Instead of $1/\theta_i$, interpolated values of $\kappa_i(x)$ which were calculated below the threshold in the two-dimensional space can also be used.

[0045] This method accordingly takes into account bleaching phenomena which vary in a spatially-dependent manner, e.g., due to different fluorophores or varying characteristics thereof.

[0046] The realization of the invention is not limited to the embodiment examples shown herein. Further developments by persons skilled in the art or other approaches for taking into account the brightness variations of the light source or of the bleaching of the object in fluorescence illumination remain within the field of the invention.